We present an on-chip wavelength reference with a partial drop ring resonator and germanium photodetector. This approach can be used in ring-resonator-based wavelength-selective switches where absolute wavelength alignment is required. We use the temperature dependence of heater resistance as a temperature sensor. Additionally, we discuss locking speed, statistical variation of heater resistances, and tuning speed of the switches. © 2017 Optical Society of America

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Wavelength-selective switches promise greater capacity and flexibility in data centers when used with electronic packet switches [1]. Multiple variants of fat pipe switches have been demonstrated in silicon photonics in [2–4]. Ring-resonator-based wavelength-selective switches have been demonstrated by different groups [5,6], and provide the advantage of fine granularity wavelength switches which better fit the traffic patterns in datacenters.

We proposed a flexible ring-resonator-based architecture in [7], where any set of multiple wavelengths could be switched between input-output pairs. In [7] we also established that two rings per cross point are sufficient for a ring-based crossbar switch where the number of wavelengths per port are equal to the number of ports. Such micro-ring-based switches that can establish a connection with any one or a set of multiple wavelengths from a WDM signal require a robust wavelength locking scheme. The locking scheme in this problem is different from other approaches [8–12], where a ring resonator or a bank of resonators can lock to a particular or all wavelengths in a WDM signal but one cannot be certain to which wavelength a particular ring is locked without an external optical spectrum analyzer (OSA). For this we need an absolute wavelength reference on chip to lock any switch node to a given wavelength of choice. In this Letter we demonstrate such an approach, which requires only an approximate tuning curve of the ring resonators, does not require an external monitor signal for wavelength identification, and can lock rings at different stage temperatures. Preliminary results from this configuration were presented in [13]. In this Letter we present locking time, statistical variation of ring resonators, and detailed discussions regarding the performance of the switch.

Figure 1(a) shows a schematic of the experimental setup of the wavelength locking experiment. A tunable laser (Yenista Tunics T100S-HP) is used to inject a wavelength at 1557 nm with TE polarization into the device. We control the polarization of light with a Newport polarization controller (PC: Model F-POL-APC). The chip is mounted on top of a thermo-electric Cooler (TEC), which we use to change the stage temperature with the help of a thermister and ILX temperature controller (Model LDT-5980). Figure 1(b) shows the experimental setup with the TEC, thermister, electrical probes, and input and outputs of the device under test (DUT).

Figure 1(c) shows the schematic of the architecture of the switch. Each switch matrix unit cell has L second-order ring resonators that can select up to L wavelengths from incoming WDM channels and a waveguide crossing. The free spectral range (FSR) of the ring resonators should be designed to fit all the channels, and the rings have to be tuned by thermal tuning as tuning across all channels is required. For example, for 8 WDM channels at 200 GHz channel spacing, an FSR of at least 1.6 THz (12.8 nm) is required. In our experiment, we use an FSR of 26 nm. The synchronous time slotted operation of the switch with centralized arbiter implies an input queued switch architecture [7]. The DUT is shown by the dashed box in Fig. 1(c) and by the micrograph of Fig. 1(d). Our locking experiment was conducted by tuning only one of the two second-order ring resonators. The optical output is recorded by an OSA (Yokogawa AQ6370).

These devices were a part of the initial release of the AIM Photonics Process Design Kit (PDK). The switch ring resonators were equipped with thermo-optic and electro-optic tuners. In our experiments we used thermal tuning due to a higher tuning efficiency of 0.84 nm/mW and negligible loss penalty as compared to electro-optic tuning. For electro-optic tuning, the resonance wavelength was blueshifted and attenuated with a loss of 3.6 dB/nm. The 3 dB bandwidth of the resonators was 40 GHz and with better than 30 dB of out of band rejection at the drop port. The measured insertion loss of each second-order
The value of resistance at this $V_{\text{WM}}$ versus the optimized wavelength $\lambda$ gives the tuning map of $R(\lambda)$ measured at stage temperature of 20°C given in Fig. 1(g). In Algorithm 1, we describe the locking procedure. Here a prior on the number of incoming WDM channels is assumed known and step 4 is useful for locking in the presence of multiple channels.

**Algorithm 1. On-Chip Locking Algorithm**

1. procedure Lock $(V_1, V_2)$ ▸ for a given $\lambda$, $R(\lambda)$
2. At $T = T_{\text{stage}}$ and $\lambda$ set $V_{\text{WM}}$ corresponding to $R(\lambda)$
3. Coarse sweep in $(V_1, V_2)$ with $V_1 = V_2$, and record PD current ($I_{PD}$)
4. Pick a guess voltage $V_{\text{guess}} = V_1 = V_2$ corresponding to peak in $I_{PD}$ corresponding to channel number
5. optimize $(V_1, V_2)$ to maximize $I_{PD}$.

Next we present our results on locking in the presence of a single channel and two channels spaced at $\Delta \nu = 100, 200$ GHz. Figure 2(a) shows the temperature dependence of the resonant wavelength. The resonant wavelength of the ring changes by 80 pm/K. In this figure, the voltages on the switch heaters were optimized for highest power transmission at 40°C and then the stage temperature was changed to 20°C and 30°C with the heaters set to voltages corresponding to 40°C. We then use our wavelength locking scheme to change the switch heater voltages with the help of the WM and lock wavelengths within 20 pm of the target wavelength. The tunable laser used for measurement had a 20 pm accuracy. The results can be seen in Fig. 2(b). The transfer function of the
The drop port spectra of the second-order ring is recorded here after the locking.

Figures 2(c) and 2(d) show before and after locking results in the presence of two channels spaced at $\Delta \nu / 0.136 \times 100 \text{GHz}$. Figures 2(e) and 2(f) show before and after locking in the presence of two channels spaced at $\Delta \nu / 0.136 \times 200 \text{GHz}$. In these experiments system loss was measured at 20 dB and input power on both wavelengths was 10 dBm. In the transfer function corresponding to $40^\circ C$ in Figs. 2(c), 2(d), 2(e), and 2(f) the peak corresponding to the other second-order switch in the unit cell with unoptimized heater voltages is also visible. The optimization is run with stochastic hill climbing (SHC) [15] with a noise standard deviation of 50 mV to limit search space around $V_{\text{guess}}$.

Figure 3(a) shows the speed of different algorithms used to optimize the heater voltages applied to the second-order switch ring. We compare three different optimization methods, namely the Nelder–Mead [16], the SHC, and Powell’s algorithms [17]. SHC, specifically in this case random mutation hill climbing, is a type of genetic algorithm where we apply stochastic noise to the voltages and update the best guess to the heater voltages depending on the fitness criteria (transmission). Because of the stochastic nature of the search, the locking time is not deterministic and this is shown by the multiple blue

![Fig. 2. Transfer function of the second-order resonator after locking to predefined wavelengths at different stage temperatures. Wavelengths to which the ring is locked is given as a black dashed line. (a) and (b) show results before and after locking in the presence of input single wavelength. (c) and (d) show results after locking to two channels separated by $\Delta \nu = 100 \text{GHz}$. (e) and (f) show results for $\Delta \nu = 200 \text{GHz}$.]

![Fig. 3. (a) Locking speed of different algorithms, with Nelder–Mead being the fastest locking algorithm. Here, the distance from optima is the separation from the terminal optimized transmission of the second-order ring resonator. (b) Measured variation of the second-order ring heater resistance across the switch chip at a stage temperature of 20°C, with an average value of 7568 Ω, and a standard deviation of 256 Ω.]

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trjectories in the figure. Both Nelder–Mead and Powell are derivative-free algorithms which give deterministic trajectories and locking times. These were implemented with the minimization module of SciPy library of Python. The top axis of locking time is generated assuming each iteration takes 7 μs, which is the measured thermal time constant. The heuristic algorithms do not terminate on a non-optimal local maxima as the second-order switch transmission is a unimodal function of heater voltages, i.e., there exists only one maximum. These heuristic algorithms are also accurate and faster than a grid search over the two heater voltages of the second-order resonators since the function is unimodal.

The main reason why we pursued this approach was that there was an 8% standard deviation in heater resistance across the switch chip, as shown in Fig. 3(b). A total of 256 heater resistances from 128 second-order rings were used to measure this standard deviation. The size of the switch die across which measurements are done is 5 mm × 2.4 mm. This means that one cannot generate a tuning curve for a single resonator and use that to lock other ring resonators as each resonator will have a different tuning curve.

In Fig. 2, we were limited by the temperature range for which our approach of absolute wavelength locking worked from 20°C to 40°C. This was due to the limited tuning of the ring resonators fabricated. Though the resonators were designed to be fully tunable across an FSR of 26 nm, we saw performance degradation of the heater after tuning by ~8 nm. Further investigation of the tuning limits of these structures is needed, but tuning across this fraction of the FSR still allows for useful switching operation.

An advantage of this method is the ability to be resistant to thermal cross talk across the switch. The wavelength monitors can be placed away from the switch rings and multiple ring resonators can be optimized simultaneously with two or more wavelength monitors in a column. For example, if two ring resonators are placed close to each other, their voltages can be optimized simultaneously with the WM as reference. These voltages will self-adjust according to the thermal cross-talk penalty to maximize drop power at a particular wavelength.

Fortunately, it was not necessary to compensate for thermal cross talk, as the observed thermal cross talk was small. In most cases, the layout of the switch is pad limited and if the pads are laid out in a flip-chip configuration, for example in a square grid at 150 μm pitch, the switch rings can be placed far apart to avoid significant thermal cross talk. This is because the footprint of a ring resonator is 20 μm × 10 μm, which is small compared to 150 μm pitch. To scale such switches, optical interposers might be used in the future and the rings can be placed closer together for tighter packing and this can increase the thermal cross talk. The motivation behind placing k WM per column in Fig. 1(c) is that speed of locking can be parallelized. We placed identical partial drop rings in our layout. If the number of channels is equal to the number of WMs then partial drop rings with different resonant wavelengths can be used.

To conclude, we present an approach to bias and stabilize ring-resonator-based wavelength-selective switches where absolute wavelength alignment is required. We use the temperature dependence of heater resistance as a temperature sensor to stabilize the ring switching wavelengths.

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